

## Transport phenomena in weak formulation

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# Advection Equation

Hyperbolic PDEs in FreeFEM++

[302.044] Numerical Methods in Fluid Dynamics

ADVECTION EQUATION

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▷ Strong form:

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▷ Strong form:

$$\begin{cases} \partial_t u - \mathbf{a} \cdot \nabla u(\mathbf{x}, t) = f(\mathbf{x}, t) & , \quad \mathbf{x} \in \Omega , t > 0 \\ +\text{Dirichlet BCs} \end{cases}$$

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$$\begin{cases} \int_{\Omega} v \partial_t u d\Omega - \int_{\Omega} v \mathbf{a} \cdot \nabla u d\Omega = \int_{\Omega} f v d\Omega & , \quad \mathbf{x} \in \Omega \\ u|_{\Gamma = \partial\Omega_{dep}} = g \end{cases}$$

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[302.044] Numerical Methods in Fluid Dynamics

CHARACTERISTIC-GALERKIN METHOD

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[302.044] Numerical Methods in Fluid Dynamics

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$$\dot{\mathbf{X}}(\tau) = \mathbf{a}(\mathbf{X}(\tau), \tau)$$

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## CHARACTERISTIC-GALERKIN METHOD

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- ▷ On the characteristic line reference frame:

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- ▷ For unforced advection ( $f = 0$ ):

$$\dot{\mathbf{X}}(\tau) = \mathbf{a}(\mathbf{X}(\tau), \tau)$$

- ▷ On the characteristic line reference frame:

$$\frac{D}{Dt}[u(\mathbf{X}(t), t)] = f(\mathbf{X}(t), t)$$

where  $\frac{D}{Dt}[\mathbf{X}(t), t] = \mathbf{a}(\mathbf{X}(t), t)$

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where  $\frac{D}{Dt}[\mathbf{X}(t), t] = \mathbf{a}(\mathbf{X}(t), t)$

- ▷ Using a 2-points FT scheme for discretizing the total derivative:

$$\frac{u^{n+1}(\mathbf{X}^{n+1}) - u^n(\mathbf{X}^n)}{\Delta t} = f^n$$

where  $\mathbf{X}^{n+1} = \mathbf{X}((k+1)\Delta t) = \mathbf{x}$

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TAYLOR EXPANSION

# Advection Equation

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TAYLOR EXPANSION

$$\begin{aligned}
 u(\mathbf{X}(k\Delta t)) &= u(\mathbf{X}((k+1)\Delta t)) - \\
 &\quad \Delta t \sum_{i=1}^{N_{dim}} \frac{\partial [u^k(\mathbf{X}((k+1)\Delta t))]}{\partial x_i} \frac{D[\mathbf{X}_i((k+1)\Delta t)]}{Dt} + o(\Delta t) \\
 &= u(\mathbf{X}((k+1)\Delta t)) - \Delta t \mathbf{a}^n(\mathbf{x}) \cdot \nabla u^n(\mathbf{x}) + o(\Delta t) \\
 &\quad \rightarrow u^n(\mathbf{x} - \mathbf{a}\Delta t) = u^n(\mathbf{x}) - \Delta t \mathbf{a}^n(\mathbf{x}) \cdot \nabla u^n(\mathbf{x}) + o(\Delta t)
 \end{aligned}$$

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 \end{aligned}$$

$$u^n(\mathbf{x} - \mathbf{a}\Delta t) \approx u^n(\mathbf{x}) - \Delta t \mathbf{a}^n(\mathbf{x}) \cdot \nabla u^n(\mathbf{x})$$

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 u(\mathbf{X}(k\Delta t)) &= u(\mathbf{X}((k+1)\Delta t)) - \\
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 &= u(\mathbf{X}((k+1)\Delta t)) - \Delta t \mathbf{a}^n(\mathbf{x}) \cdot \nabla u^n(\mathbf{x}) + o(\Delta t) \\
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 \end{aligned}$$

$$u^n(\mathbf{x} - \mathbf{a}\Delta t) \approx u^n(\mathbf{x}) - \Delta t \mathbf{a}^n(\mathbf{x}) \cdot \nabla u^n(\mathbf{x})$$

▷ Advection Equation: FreeFEM++ `convect` operator:

$$\text{convect}(\mathbf{a}, -\Delta t, u^n) = u^n(\mathbf{x}) - \Delta t \mathbf{a}^n(\mathbf{x}) \cdot \nabla u^n(\mathbf{x}) \approx u^{n+1}(\mathbf{x})$$



# Advection Equation

[302.044] Numerical Methods in Fluid Dynamics

Hyperbolic PDEs in FreeFEM++

## 2-D ADVECTION EQUATION

▷ First case:  $\Omega = [0, 1] \times [0, 1]$  and  $\mathbf{a}$  constant

```

1  mesh Omega = square(60,20);
2  fespace Vh(Omega,P1);
3  Vh u0 = exp(-10*((x-0.3)^2 +(y-0.3)^2));
4  Vh a1 = 1, a2 = 1; u;
5  real dt = 0.0017,t=0;
6  for (int i=0; i<200; i++) {
7    t += dt;
8    u = convect([a1,a2],-dt,u0);
9    u0 = u;
10 plot(u,fill=1,wait=0);
11 };

```

## Advection Equation

## Object Oriented Programming in FreeFEM++

## FREEFEM++: OBJECT ORIENTED PROGRAMMING

- ▷ First case:  $\Omega = [0, 1] \times [0, 1]$  and **a** constant

## Mesh Block

```
1 mesh Omega = square(60,20);
```

## Functional Space Block

```
1 fespace Vh(Omega,P1);
2 Vh u0 = exp(-10*((x-0.3)^2 + (y-0.3)^2));
3 Vh a1 = 1, a2 = 1; u;
```

## Solver Block

```
1 real dt = 0.0017, t=0;
2 for (int i=0; i<200; i++) {
3 t += dt; u = convect([a1,a2],-dt,u0);
4 u0 = u; plot(u,fill=1,wait=0); }
```

## Advection Equation

## Object Oriented Programming in FreeFEM++

## FREEFEM++: OBJECT ORIENTED PROGRAMMING

▷ Second case:  $\Omega = \{x^2 + y^2 \leq 1\}$  and a variable

```

1  border C(t=0, 2*pi) {x=cos(t); y=sin(t)};
2  mesh Omega = buildmesh(C(200));
3  fespace Vh(Omega, P1);
4  Vh u0 = exp(-10*((x-0.3)^2 + (y-0.3)^2));
5  Vh a1 = y, a2 = -x; u;
6  real dt = 0.17, t=0;
7  for (int i=0; i<200; i++) {
8  t += dt;
9  u = convect([a1, a2], -dt, u0);
10 u0 = u;
11 plot(u, fill=1, wait=0);
12 };

```

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▷ Second case:  $\Omega = \{x^2 + y^2 \leq 1\}$  and a variable

Mesh Block

Changed

```
1 border C(t=0,2*pi) {x=cos(t); y=sin(t)};
2 mesh Omega = buildmesh(C(200));
```

Functional Space Block: Changed

```
1 fespace Vh(Omega,P1);
2 Vh u0 = exp(-10*((x-0.3)^2 +(y-0.3)^2));
3 Vh a1 = y, a2 = -x; u;
```

Solver Block:

Unchanged

# Wave Equation

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$$\begin{cases} \partial_{tt}p - c^2 \nabla^2 p(\mathbf{x}, t) = f(\mathbf{x}, t) & , \quad \mathbf{x} \in \Omega , t > 0 \\ + \text{Dirichlet BCs} \end{cases}$$

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▷ Weak form:

$$\begin{cases} \int_{\Omega} v \partial_{tt} p d\Omega + \int_{\Omega} c^2 \nabla p \cdot \nabla v d\Omega = \int_{\Omega} f v d\Omega & , \quad \mathbf{x} \in \Omega \\ p|_{\Gamma = \partial\Omega_{dep}} = g \end{cases}$$

## Wave Equation

## Hyperbolic PDEs in FreeFEM++

## 2-D WAVE EQUATION

▷ First case:  $\Omega = \{x^2 + y^2 \leq 1\}$  constant

```

1  border Gamma(t=0, 2*pi) { x=cos(t); y=sin(t); };
2  mesh Omega = buildmesh(Gamma(200));
3  fespace Vh(Omega, P2);
4  Vh u, uh;
5  Vh uprev = 5*exp(-100*((x)^2+(y)^2));
6  Vh uprev2 = 5*exp(-100*((x)^2+(y)^2));
7  real dt = 0.01, t=0;
8  macro Grad(u) [dx(u), dy(u)]
9  problem solveWave(u, uh)
10 = int2d(Omega)(2*Grad(u){'} * Grad(uh) + u*uh/dt/dt)
11 + int2d(Omega)((-2*uprev + uprev2) * uh/dt/dt)
12 + on(1,2,3,4,u=0);

```

## Wave Equation

## Hyperbolic PDEs in FreeFEM++

## 2-D WAVE EQUATION

▷ First case:  $\Omega = \{x^2 + y^2 \leq 1\}$  constant

```

1  real[int] viso(60);
2  for(int i=0;i<viso.n;i++)
3      viso[i]=i*0.05-1.1;
4  for (int i=0; i<5000; i++)
5  {
6      t += dt; solveWave; uprev2 = uprev; uprev = u;
7      if (i%10==0)
8          {
9              plot(u,fill=1,wait=0,viso=viso(0:viso.n-1),dim=3,
10                 value=true,cmm="t = "+t+" i "+i );
11          }
12  };

```

# Isentropic Inviscid Flow

Hyperbolic PDEs in FreeFEM++

[302.044] Numerical Methods in Fluid Dynamics

ISENTROPIC INVISCID FLOW<sup>1</sup>

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<sup>1</sup>Under the hypothesis of quasi-stationarity near the boundaries.  
O. Pironneau, *Corrected Nonconservative Scheme*,  
Chinese Annals of Mathematics (1/6/5)

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# Isentropic Inviscid Flow

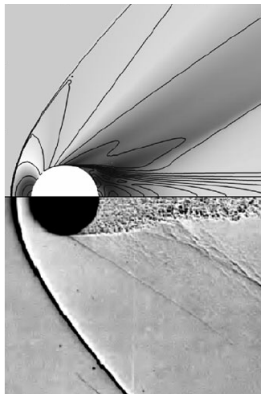
Hyperbolic PDEs in FreeFEM++

[302.044] Numerical Methods in Fluid Dynamics

ISENTROPIC INVISCID FLOW<sup>1</sup>

▷ Strong form:

$$\left\{ \begin{array}{l} \rho_{,t} + \nabla \cdot (\rho \mathbf{u}) = 0 \\ (\rho \mathbf{u})_{,t} + \nabla \cdot (\rho \mathbf{u} \otimes \mathbf{u} + \rho \gamma) = 0 \\ + \text{Dirichlet BCs} \end{array} \right.$$



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## Isentropic Inviscid Flow

Hyperbolic PDEs in FreeFEM++

## ISENTROPIC INVISCID FLOW

▷ Strong form:

$$\begin{cases} \frac{1}{\bar{\rho}^n \Delta t} (\rho^{n+1} - \rho^n \circ X^n) + \nabla \cdot \mathbf{u}^{n+1} = 0 \\ \frac{\rho^{n+1} - \rho^n}{\gamma \Delta t} (\mathbf{u}^{n+1} - \mathbf{u}^n \circ \tilde{X}^n) + \nabla \rho^{n+1} = 0 \end{cases}$$

where  $\tilde{X}$  is computed with the velocity  $\overline{\rho \mathbf{u}} / \bar{\rho}$ .

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$$\begin{cases} \frac{1}{\bar{\rho}^n \Delta t} (\rho^{n+1} - \rho^n \circ X^n) + \nabla \cdot \mathbf{u}^{n+1} = 0 \\ \frac{\rho^{n+1} - \rho^n}{\gamma \Delta t} (\mathbf{u}^{n+1} - \mathbf{u}^n \circ \tilde{X}^n) + \nabla \rho^{n+1} = 0 \end{cases}$$

where  $\tilde{X}$  is computed with the velocity  $\overline{\rho \mathbf{u}} / \bar{\rho}$ .

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$$\begin{cases} \frac{1}{\overline{\rho^n} \Delta t} (\rho^{n+1} - \rho^n \circ X^n) + \nabla \cdot \mathbf{u}^{n+1} = 0 \\ \frac{\overline{\rho^{n^2-\gamma}}}{\gamma \Delta t} (\mathbf{u}^{n+1} - \mathbf{u}^n \circ \tilde{X}^n) + \nabla \rho^{n+1} = 0 \end{cases}$$

where  $\tilde{X}$  is computed with the velocity  $\overline{\rho \mathbf{u}} / \bar{\rho}$ .

▷ Weak form:

$$\begin{cases} \int_{\Omega} \frac{w_h \rho_h^{n+1}}{\rho_h^n} - \Delta t \left( \int_{\Omega} \mathbf{u}_h^{n+1} \cdot \nabla w_h - \int_{\Gamma} \mathbf{u}_h^{n+1} \cdot \mathbf{n} w_h \right) = \int_{\Omega} \frac{\rho_h^n \circ X^n w_h}{\rho_h^n} \\ \int_{\Omega} \overline{\rho_h^{n^2-\gamma}} \mathbf{u}_h^{n+1} \mathbf{v}_h + \gamma \Delta t \int_{\Omega} \nabla \cdot (\rho^{n+1} \mathbf{v}_h) = \int_{\Omega} \overline{\rho^{n^2-\gamma}} \mathbf{u}_h^n \circ X^n \mathbf{v}_h \end{cases}$$

## Isentropic Inviscid Flow

[302.044] Numerical Methods in Fluid Dynamics

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## ISENTROPIC INVISCID FLOW

```

1  real pi2=atan(1.0)*2, R=0.3, x1=-3, x2=1.0, y2=1.5;
2  border a1(t=x1,-2-R){x=t; y=0;}
3  border a2(t=-2+R,x2){x=t; y=0;}
4  border a3(t=0,y2){x=x2;y=t;}
5  border a4(t=x2,x1){x=t;y=y2;}
6  border a5(t=y2,0){x=x1;y=t;}
7  border bb(t=pi2*2,0){x=-2+R*cos(t);
8                          y=R*sin(t);}
9  mesh Th = buildmesh(a1(30)+a2(40)+a3(20)+a4(60)+a5(20)+
10                      bb(30));
11 plot(Th,wait=1);
12 real dt=0.02, u0=2*sqrt(1.4),
    v0=0, visc=0.00125;

```

## Isentropic Inviscid Flow

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## ISENTROPIC INVISCID FLOW

```

1  fespace Vh(Th,P1);
2  Vh u,v,r,r1,u1,v1,rh,uh,vh,fro;
3  problem eul(u,v,r,uh,vh,rh)
4  = int2d(Th)((fro*(u*uh+v*vh)+r*rh/r1)/dt
5  + (dx(r)*uh+ dy(r)*vh-dx(rh)*u-dy(rh)*v))
6  + int2d(Th)(-(rh*convect([u1,v1],-dt,r1)/r1
7  + fro*(uh*convect([u1,v1],-dt,u1)
8  + vh*convect([u1,v1],-dt,v1)))/dt)
9  + int1d(Th,3)(rh*u) + int1d(Th,4)(rh*v)
10 + on(5,r=1) + on(5,u=u0) + on(1,2,5,v=0);
11 u1= u0; v1= v0; r1 = 1;
12 for(int k=0;k<200;k++){ // the time loop
13 fro = pow(r1,0.6)/1.4; eul;
14 u1=u;r1=abs(r);
15 v1=v*(y<1.4)+abs(v)*(y>=1.4);
16 // to avoid reflexion from top
17 plot(r,wait=0,value=1);
18 }

```